Tailoring light propagation through controllable phase and group velocities

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The visuals of this talk are posted at boydnlo.ca/presentations

Tailoring light propagation through controllable phase and group velocities

- Controllable group velocity
  "Slow" and "fast" light
  Leads to many effects including light-drag

- Controllable phase velocity
  Epsilon-near-zero (ENZ) materials

Controlling the Velocity of Light

"Slow," "Fast" and "Backwards" Light

– Light can be made to go:

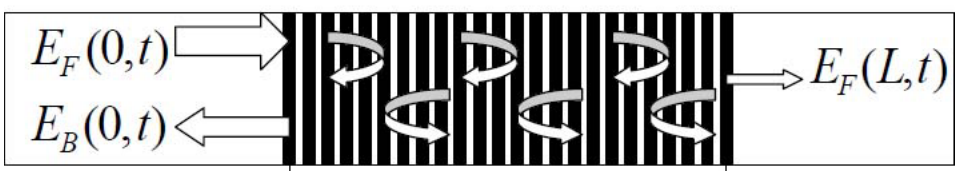
slow: $v_g << c$ (as much as 10^6 times slower!)

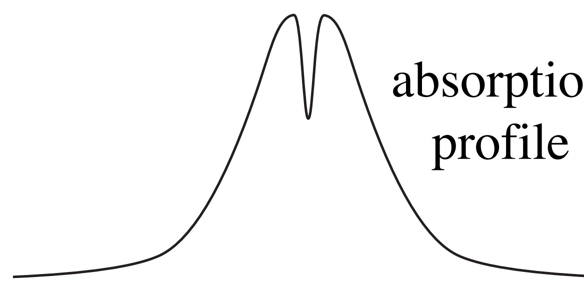
fast: $v_g > c$

backwards: v_g negative

Here v_g is the group velocity: $v_g = c/n_g \quad n_g = n + \omega\,(dn/d\omega)$

– Velocity controlled by structural or material resonances

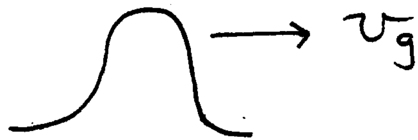




Review article: Boyd and Gauthier, Science 326, 1074 (2009).

Group Velocity

Pulse
(wave packet)



Group velocity given by $v_g = \frac{d\omega}{dk}$

For $k = \frac{n\omega}{c}$. $\frac{dk}{d\omega} = \frac{1}{c}\left(n + \omega \frac{dn}{d\omega}\right)$

Thus

$$v_g = \frac{c}{n + \omega \frac{dn}{d\omega}} \equiv \frac{c}{n_g}$$

Thus $n_g \neq n$ in a dispersive medium!

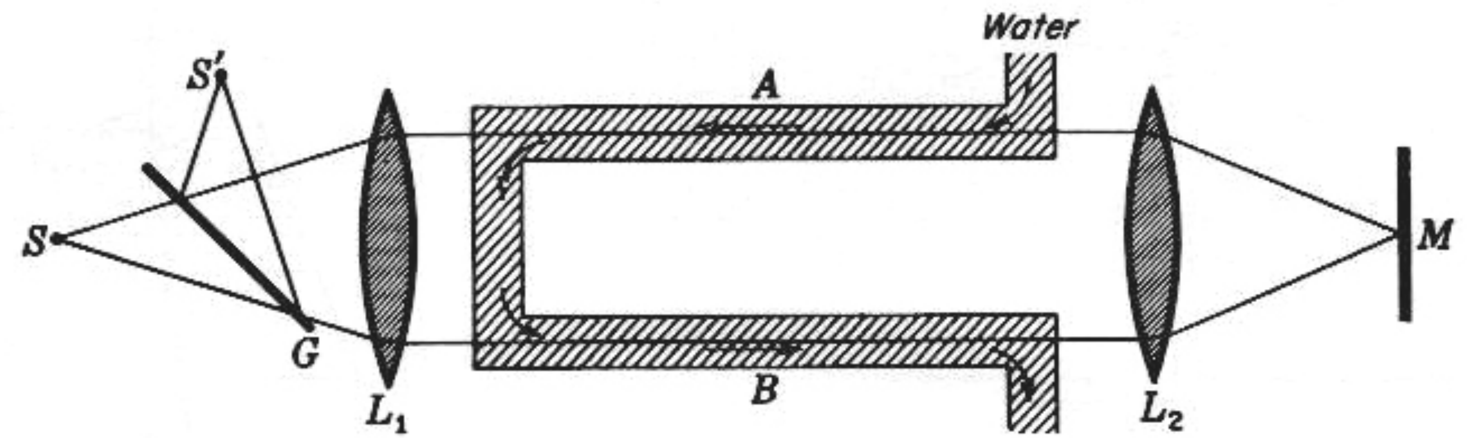
Photon Drag and Slow Light

The Velocity of Light in Moving Matter: Fresnel Drag (or Ether Drag) Effects

- Fizeau (1859): Longitudinal photon drag:

  Velocity of light in flowing water.

  V = 700 cm/sec; L = 150 cm; displacement of 0.5 fringe.



- Modern theory: relativistic addition of velocities

$$v = \frac{c/n + V}{1 + (V/c)(1/n)} \approx \frac{c}{n} + V\left(1 - \frac{1}{n^2}\right)$$

Fresnel "drag" coefficient

- But what about slow-light media?

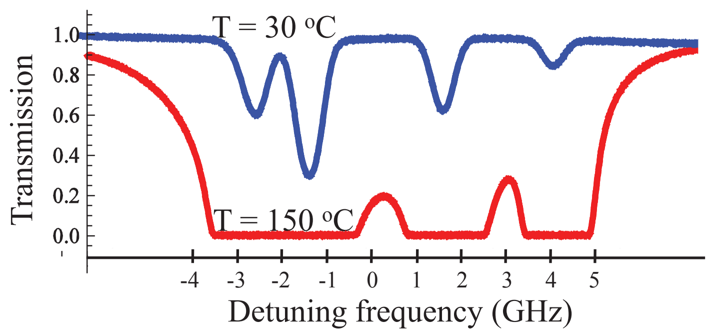
Fresnel Drag in a Highly Dispersive Medium

Light Drag in a Slow Light Medium (Lorentz)

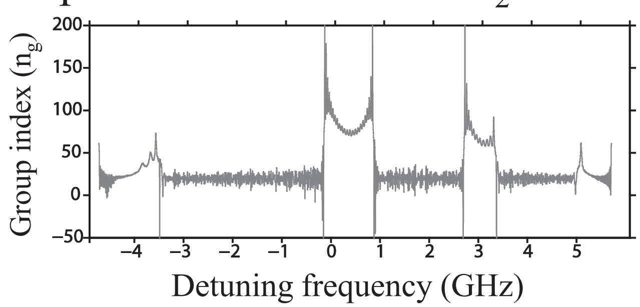
$$u \simeq \frac{c}{n} \pm v\left(1 - \frac{1}{n^2} + \frac{n_g - n}{n^2}\right)$$

We Use Rubidium as Our Slow Light Medium

- Transmission spectrum of Rb around D_2 transition:

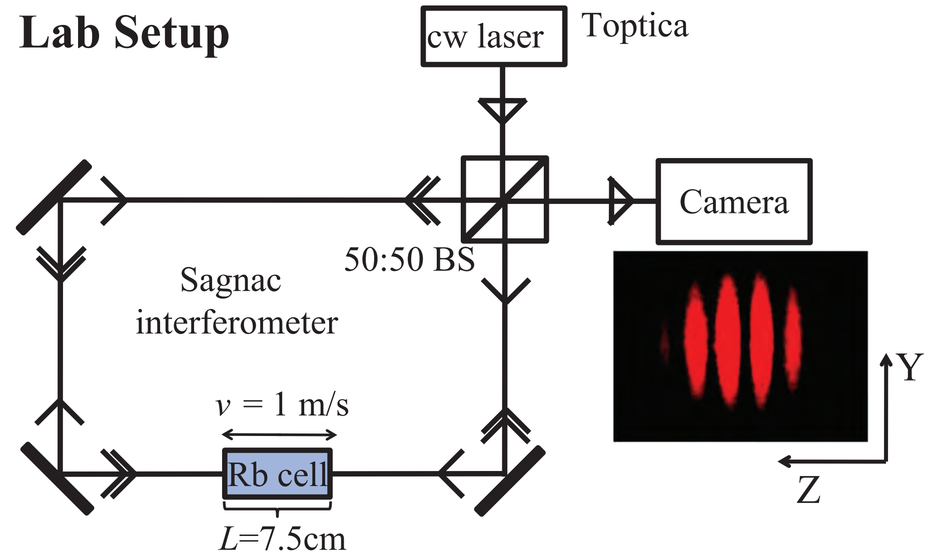


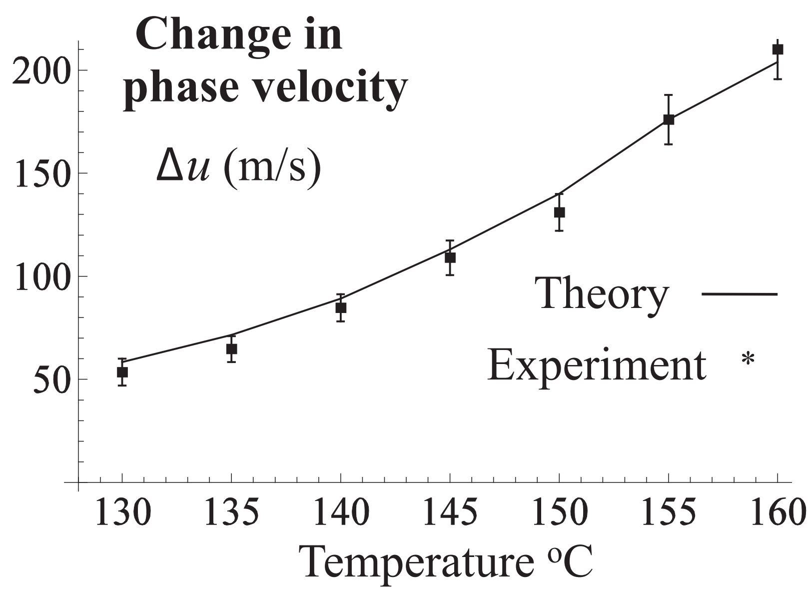
- Group index of Rb around D_2 line at T=130



Safari, De Leon, Mirhosseini, Magana-Loaiza, and Boyd
Phys. Rev. Lett. 116, 013601 (2016)

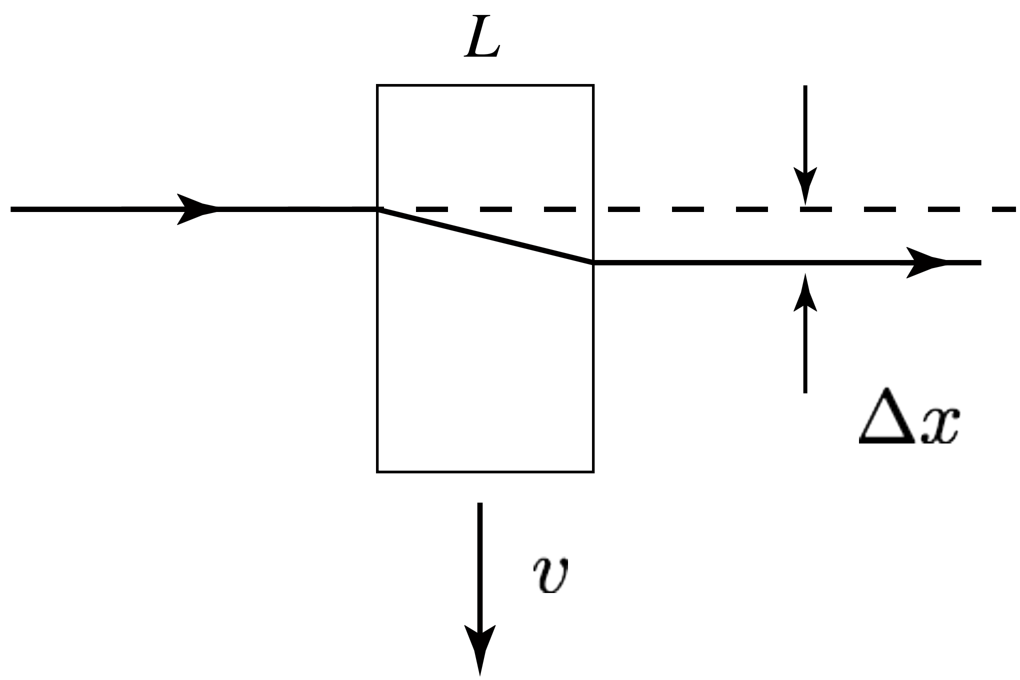
Lab Setup





- **Change in phase velocity is much larger than velocity of rubidium cell. Implications for new velocimeters?**

Transverse Photon Drag

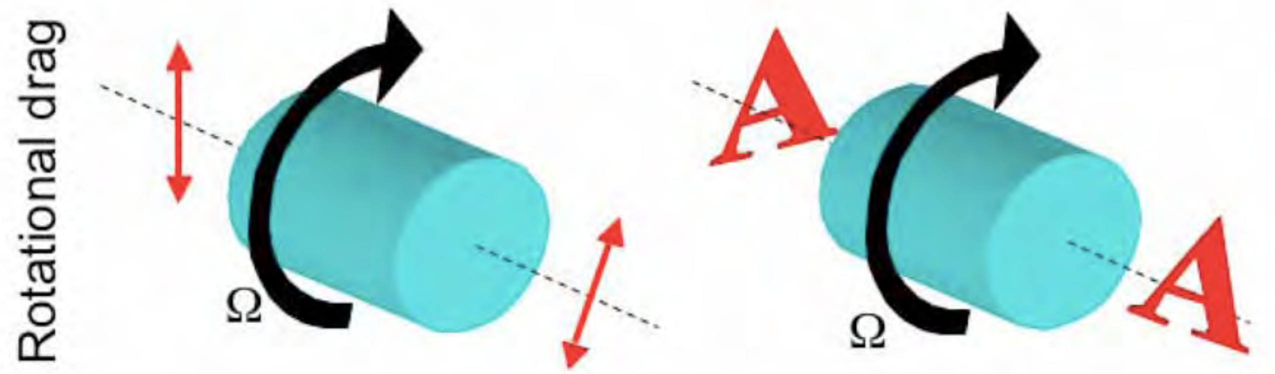


$$\Delta x = (vL/c)(n_g - 1/n_\phi)$$

For L = 25 mm, v = 2000 cm/s, displacement = 6 nm.

Measured by R.V. Jones, 1972.

Rotary Photon Drag:
An image viewed through a spinning window



$$\Delta\theta_{pol} = \Omega\left(n_g - 1/n_\phi\right)L/c \qquad \Delta\theta_{image} = \Omega\left(n_g - 1/n_\phi\right)L/c$$

Theory says that transmitted image is rotated! (rotary photon drag)

(Polarization rotation measured earlier by Jones.)

Image rotation never previously observed
(although implied by work of Leach et al., PRL 2008.)

Effect scales as group index!

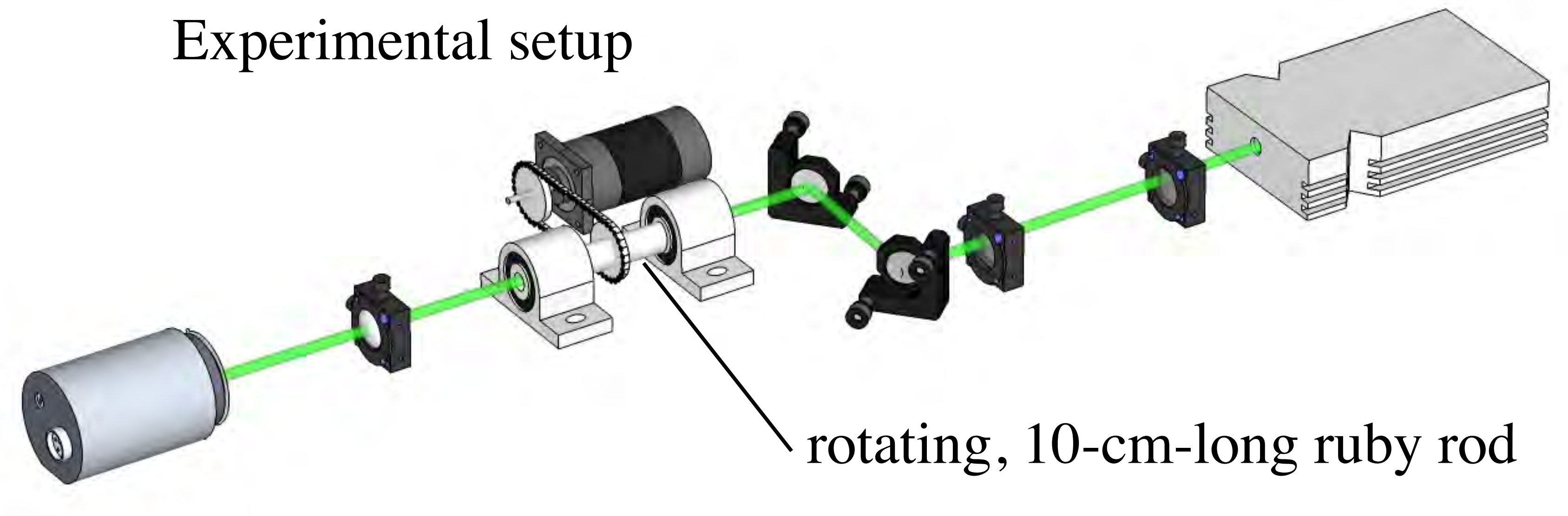
Franke-Arnold, Gibson, Boyd and Padgett, Science, 2011

Observation of Rotary Photon Drag

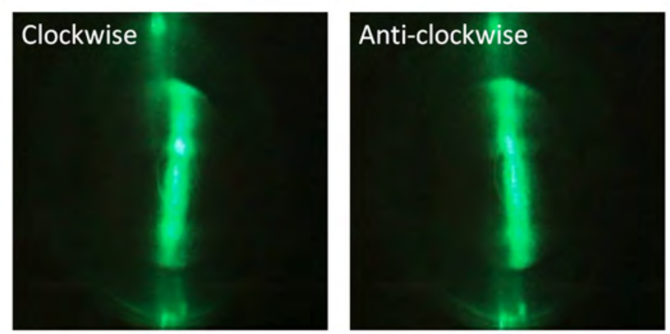
The world as seen through a spinning window.

(Laser-excited ruby has a group index of 10^6.)

Experimental setup



rotating, 10-cm-long ruby rod



Effect clearly visible by eye!

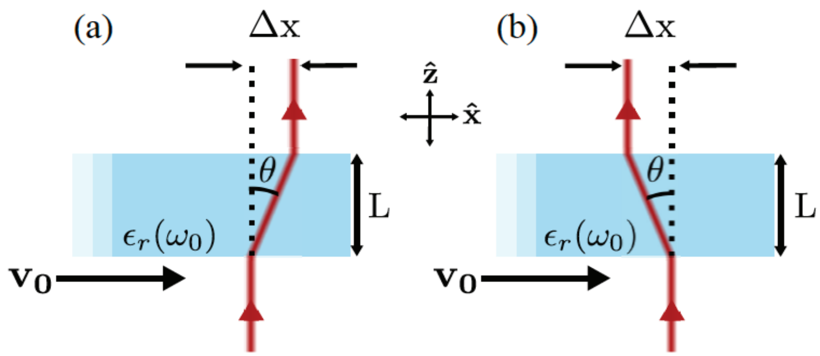
Franke-Arnold, Gibson, Boyd and Padget, Science, 2011

Anomalous Optical Drag

Chitram Banerjee,[1] Yakov Solomons,[1] A. Nicholas Black,[2] Giulia Marcucci,[3]
David Eger,[1] Nir Davidson,[1] Ofer Firstenberg,[1] and Robert W. Boyd[4,3]

(a) Normal drag; n_g is positive, beam dragged in same direction as v_0

(b) anomalous drag; n_g is negative, beam dragged in opposite direction as v_0

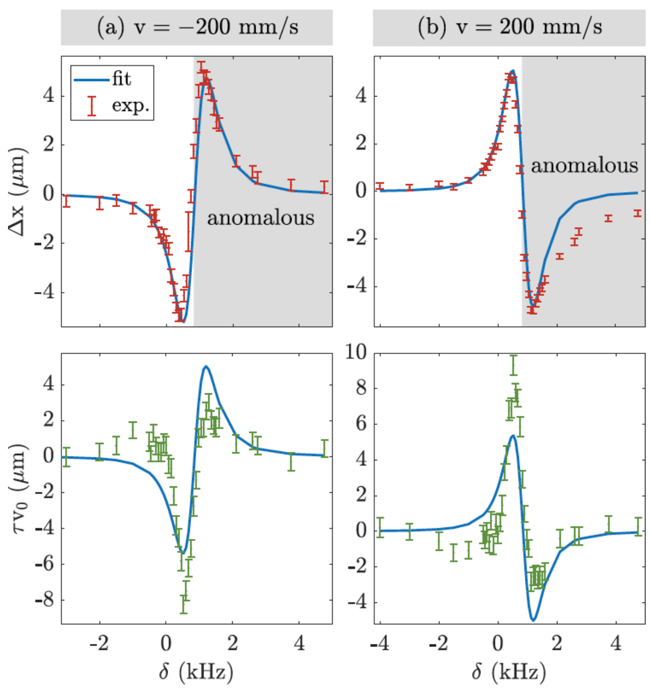


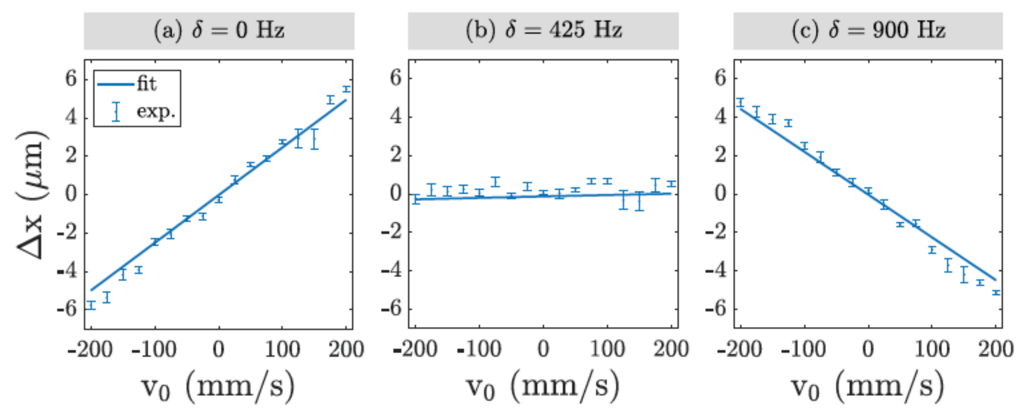
$$\Delta x = L \tan\theta. \quad (1)$$

In Eq. (1), L is the medium's longitudinal length, and θ the light's walk-off angle inside the medium, with

$$\tan\theta = \frac{v_0}{c}\left(\frac{c}{v_g} - \frac{v_p}{c}\right), \quad (2)$$

Perform experiment in Rb vapor under EIT conditions.





(a) positive group velocity ($n_g = 1.6 \times 10^7$)
(b) zero group velocity
(c) negative group velocity ($n_g = -8 \times 10^5$)

Anomalous Optical Drag

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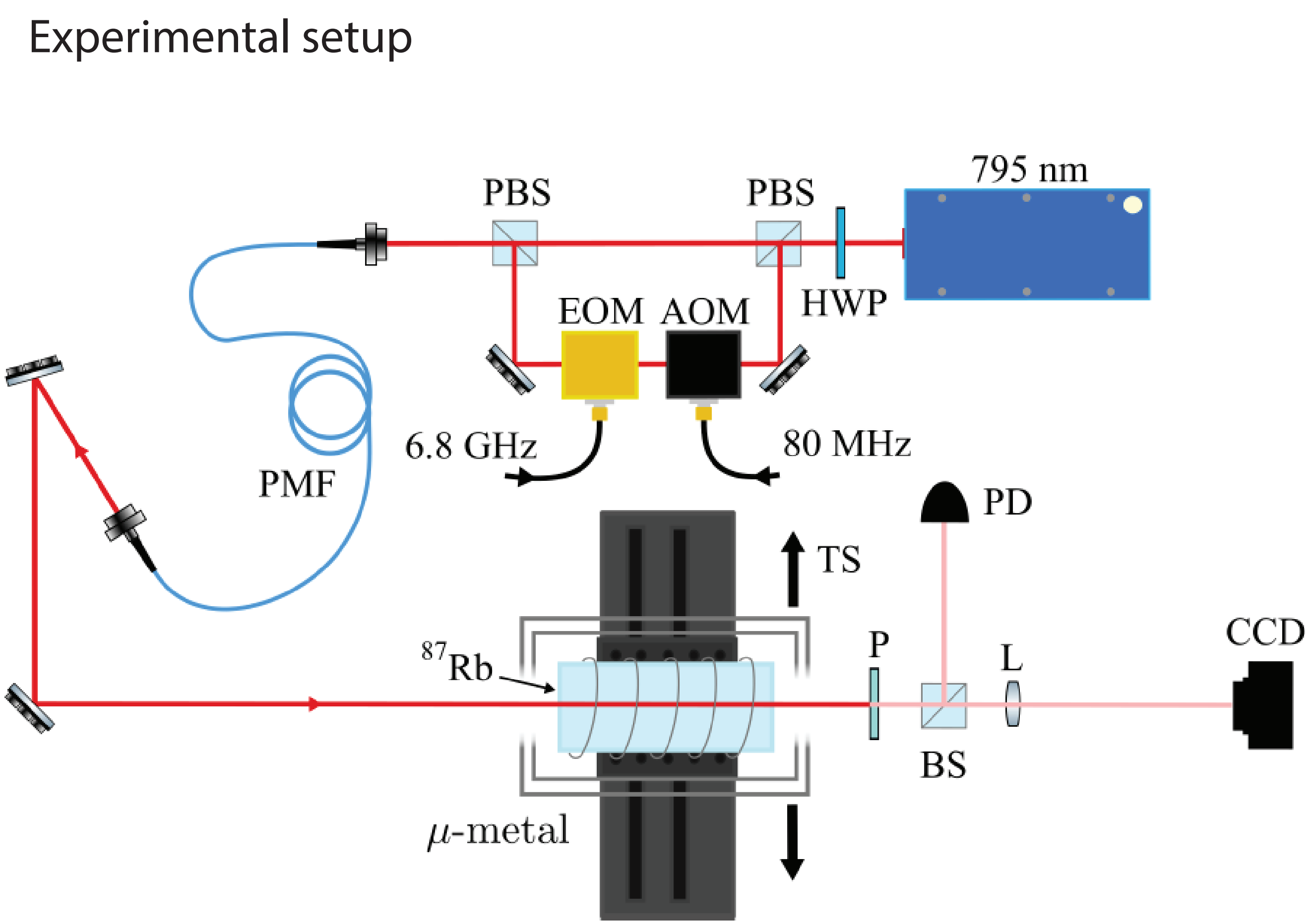
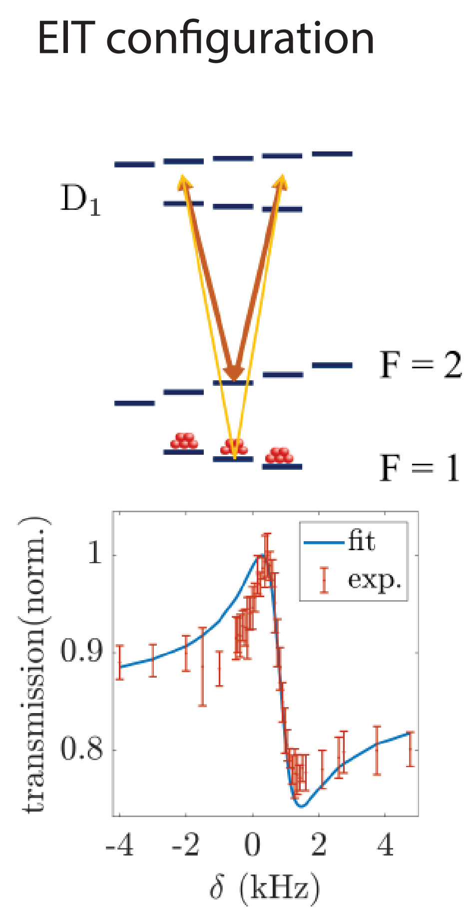




Tailoring light propagation through controllable phase and group velocities

- Controllable group velocity
  "Slow" and "fast" light
  Leads to many effects including light-drag

- Controllable phase velocity
  Epsilon-near-zero (ENZ) materials

How Light Behaves when the Refractive Index Vanishes

- Physics of Near-Zero Index (NZI) and Epsilon-Near Zero (ENZ) Materials
- Nonlinear Optical Properties of ENZ and NZI Materials
- Metamaterials for ENZ and NZI Studies
- Applications of ENZ and NZI Materials

Spoiler Alert: Implications of ENZ Behavior for Nonlinear Optics

Here is the intuition for why the ENZ condition is of interest in NLO

Recall the standard relation between n_2 and $\chi^{(3)}$

$$n_2 = \frac{3\chi^{(3)}}{4\epsilon_0 c\, n_0 \,\mathrm{Re}(n_0)}$$

Note that under ENZ conditions the denominator becomes very small, leading to a very large value of n_2

Footnote:

Standard notation for perturbative NLO

$$P = \chi^{(1)}E + \chi^{(2)}E^2 + \chi^{(3)}E^3 + \ldots$$

P is the induced dipole moment per unit volume and E is the field amplitude.

Also, the refractive index changes according to

$$n = n_0 + n_2 I + n_4 I^2 + \ldots$$

Giant Nonlinear Response of ENZ Metastructures

- Nonlinear Optics is important for a variety of reasons:

  Photonic Devices

  All-optical switching, buffers and routers based on slow light

  Used to create quantum states of light for

  Quantum Computing/Communications/Imaging

  Fundamental understanding of light-matter interactions

  Not "just" Lorentz oscillator formalism

  Understand rogue waves

  Induce and control filamentation processes

- However, the nonlinear response is usually much weaker than the linear response

- Means to enhance the nonlinear response

  Resonance interactions (atomic vapors)

  Plasmonic systems

  Electromagnetically induced transparency (EIT)

  Metamaterials (composite materials)

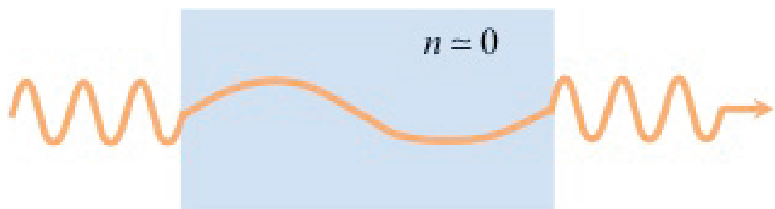
- Our approach: Use epsilon-near-zero (ENZ) materials and metamaterials

Physics of Near-Zero-Index (NZI) and Epsilon-Near-Zero (ENZ) Materials

- The wavelength of light is given by

$$\lambda = \lambda_{\text{vac}}/n$$

and is significantly lenthened in a NZI material. The wavelength approaches infinity as *n* approaches zero.



- The phase velocity of light is given by

$$v = c/n$$

and also approaches infinity as n approaches zero.

- For n approaching zero, the field oscillates in time but not in space; oscillations are in phase everywhere

Brown, Proc. IEE 100, 5 (1953).
Ziolkowski, Phys. Rev. E 70, 046608 (2004).
Silveirinha and Engheta, Phys. Rev. Lett. 97, 157403 (2006).

Physics of Epsilon-Near-Zero (ENZ) Materials

- Radiative processes are modified in ENZ materials *

Einstein A coefficient (spontaneous emission lifetime = $1/A$)

$A = n A_{\text{vac}}$

We can control (inhibit!) spontaneous emission!

Einstein B coefficient

Stimulated emission rate = B times EM field energy density

$B = B_{\text{vac}} / (n n_{\text{g}})$

Optical gain is very large!

Einstein, Physikalische Zeitschrift 18, 121 (1917).
Milonni, Journal of Modern Optics 42, 1991 (1995).

Equations are shown for nonmagnetic ($\mu = 1$) materials

- Implications:

- If we can inhibit spontaneous emission, we can build thresholdless lasers.
- Expect superradiance effects to be pronounced in ENZ materials.

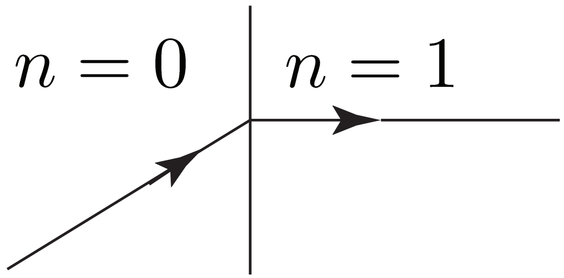
* Lobet, Liberal, Knall, Alam, Reshef, Boyd, Engheta, and Mazur, ACS Photonics 7, 1965-1970 (2020).

Optics of Zero-Index Materials

- Snell's law leads to intriguing predictions

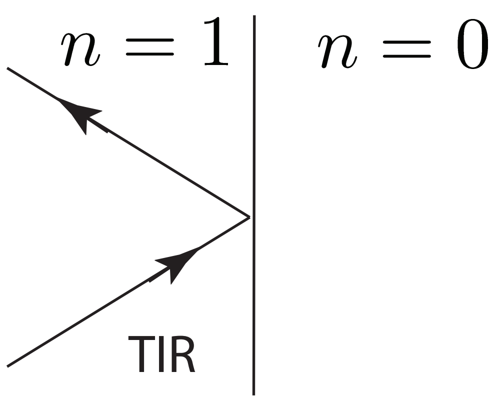
$$n_1 \sin \theta_1 = n_2 \sin \theta_2$$

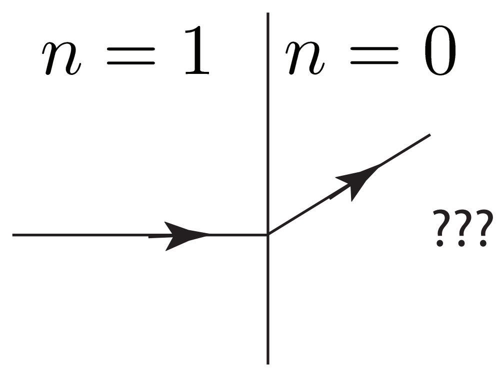
- Light always leaves perpendicular to surface of ENZ material!

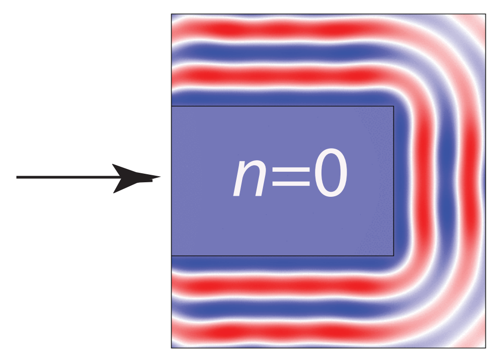


Y. Li, et al., Nat. Photonics 9, 738, 2015; D. I. Vulis, et al., Opt. Express 25, 12381, 2017.

- Thus light can enter an ENZ material only at normal incidence!







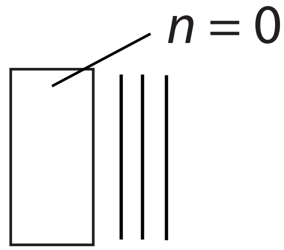
Light enters at normal incidence but leaves in all directions.

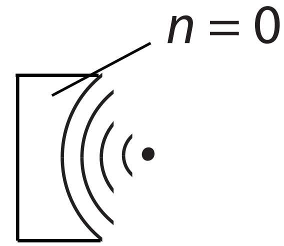
Y. Li, et al., Nat. Photonics 9, 738, 2015.

(wave-optics simulation - O. Reshef)

Some Consequences of ENZ Behavior - 1

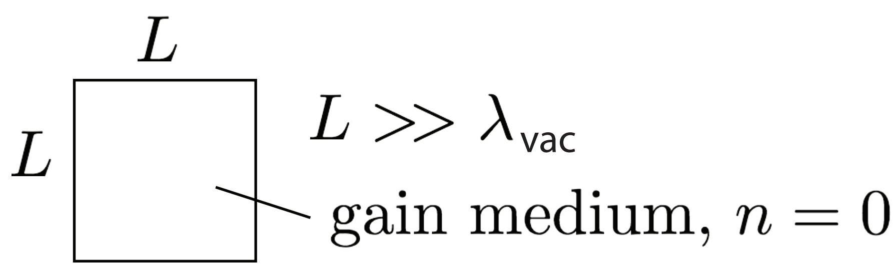
- “Funny” lenses





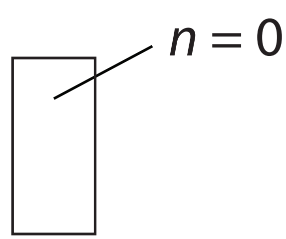
A. Alù et al., Phys. Rev. B 75, 155410, 2007; X.-T. He, ACS Photonics, 3, 2262, 2016.

- Large-area single-transverse-mode surface-emitting lasers



J. Bravo-Abad et al., Proc. Natl. Acad. Sci. USA 109, 976, 2012.

- No Fabry-Perot interference



O. Reshef et al., ACS Photonics 4, 2385, 2017.

Some Details from Electromagnetic Theory

- The linear response of any material to electromagnetic radation can be described by

  - The dielectric permittivity (dielectric constant) ε define through the relation

$$\mathbf{D} = \epsilon_0 \epsilon \mathbf{E}$$

where **D**, known as the dielectric displacement, and **E**, known as the electric field, are the two fieldsthat describe the material response to an electric field.

  - The magnetic permeability μ define through the relation

$$\mathbf{B} = \mu_0 \mu \mathbf{H}$$

where **B,** known as the magnetic field, and **H**, known as the magnetic intensity, are the two fields that describe the magnet response of a material to an applied field.

- It is straightforward to shown from the equations of electromagnetism that

$$n = \sqrt{\epsilon \mu}$$

- Thus, n=0 when either ε =0 or μ=0 (or both ε and μ equal zero).

- Terminology:

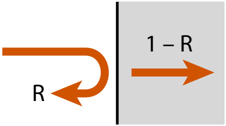
  ENZ: epsilon near zero
  MNZ: mu near zero
  EMNZ: epsilon and mu near zero

Surface Reflection

- There is a problem getting light into a zero-index material.
- There is always reflection from the boundary between two materials
- The impedance and surface reflectivity are given by

$$Z = \sqrt{\mu/\epsilon} \qquad R = \left|\frac{Z-1}{Z+1}\right|^2$$

- Thus the reflectivity will be 100% if $\epsilon = 0$ unless $\mu = 0$ as well.



- This is one reason for the interest in developing EMNZ materials (epsilon and mu near zero materials.

How Light Behaves when the Refractive Index Vanishes

- Physics of Near-Zero Index (NZI) Materials
- Nonlinear Optical Properties of NZI Materials
- Meta-materials for NZI Studies
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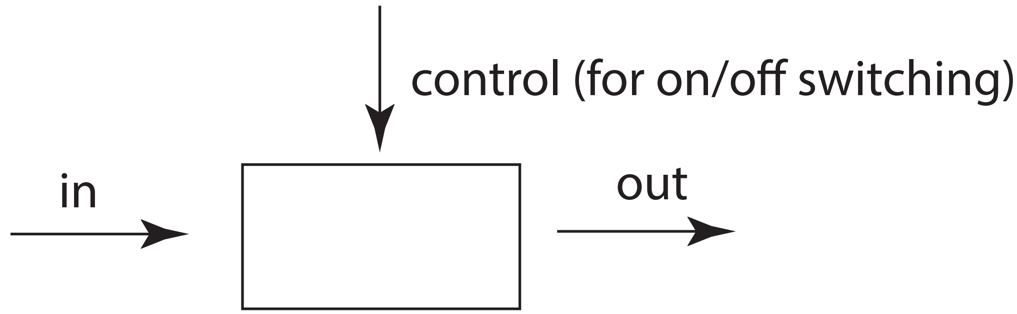
Nonlinear Optics and Optical Switching

- An important application in photonic technologies is optical switching.



- One wants a switch with fast switching times and that operates with weak control fields.
- One needs a nonlinear interaction in order for one optical field to control another field.
- A strong nonlinear response is needed. How does one quantify the strength of a nonlinear response? Two standard methods:

$$n = n_0 + n_2 I$$

$$P^{\mathrm{NL}} = 3\chi^{(3)}|E|^2 E$$

- The nonlinear coefficients are n_2 and $\chi^{(3)}$

Implications of ENZ Behavior for Nonlinear Optics

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Footnote:

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$$P = \chi^{(1)}E + \chi^{(2)}E^2 + \chi^{(3)}E^3 + \ldots$$

P is the induced dipole moment per unit volume and E is the field amplitude.

Also, the refractive index changes according to

$$n = n_0 + n_2 I + n_4 I^2 + \ldots$$

How to Choose an Epsilon-Near-Zero Materials

- Electrical conductors

  All conductors display ENZ behavior at their (reduced) plasma frequency

  Recall the Drude formula

$$\epsilon(\omega) = \epsilon_\infty - \frac{\omega_p^2}{\omega(\omega + i\gamma)}$$

  Note that $\mathrm{Re}\,\epsilon = 0$ for $\omega = \omega_p/\sqrt{\epsilon_\infty} \equiv \omega_0$.

  ENZ wavelength restricted to a limited range in the visible.

- Electrical insulators (dielectrics)

  Dielectrics can show ENZ behavior at their (optical) phonon resonance.

  ENZ wavelength restricted to a limited range in the mid-IR.

- Metamaterials

  Can design the material so that the ENZ or EMNZ wavelengths are at any desired value.

- Challenge (for any material system). For low loss, we want Im ε as small as possible at the wavelength where Re ε =0.

Nonlinear Optics of Indium Tin Oxide (ITO)

- We recently reported that, at its ENZ wavelength, ITO possesses a nonlinear coefficient n_2 that is 100 times larger than those of previously reported materials [1].
- ITO is a degenerate semiconductor (so highly doped as to be metal-like).
- ITO has a large density of free electrons, and a bulk plasma frequency corresponding to a wavelength of approximately 1.24 μm.
- Dielectric properties of ITO are well described by the Drude formula.

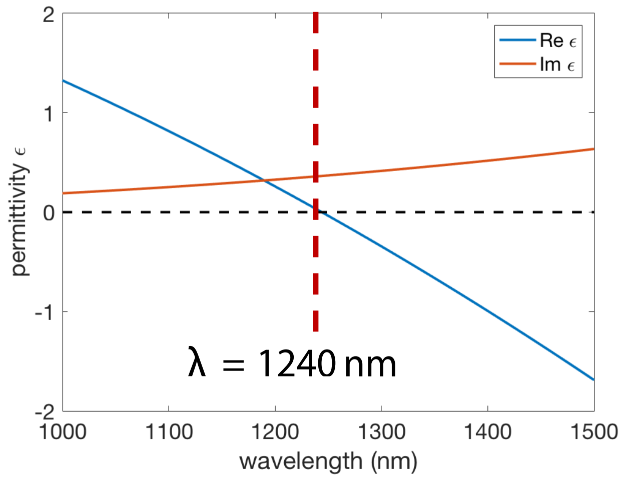
$$\epsilon(\omega) = \epsilon_\infty - \frac{\omega_p^2}{\omega(\omega + i\gamma)}$$

- Note that aluminum-doped zinc oxide (AZO), another transparent conducting oxide, also has strong nonlinear response at its ENZ wavelength [2].

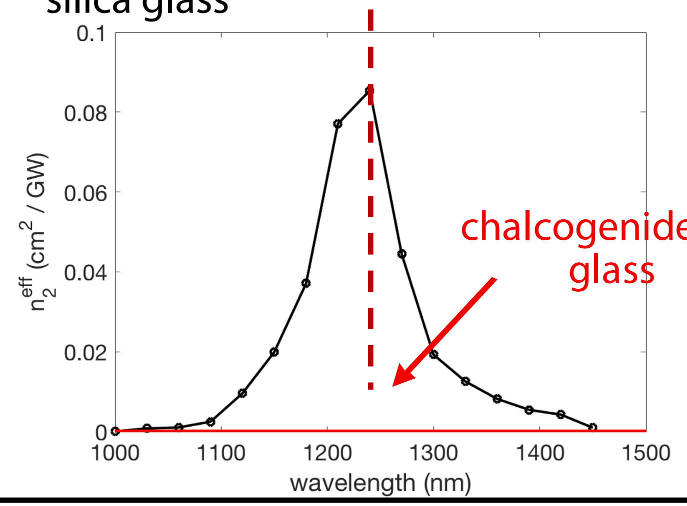
1. Alam, De Leon and Boyd, Science 352, 795–797 (2016)
2. Caspani, Shalaev, Boltasseva, Faccio et al., Phys. Rev. Lett. 116, 233901 (2016).

Huge, Fast NLO Response of Indium Tin Oxide at its ENZ Wavelength

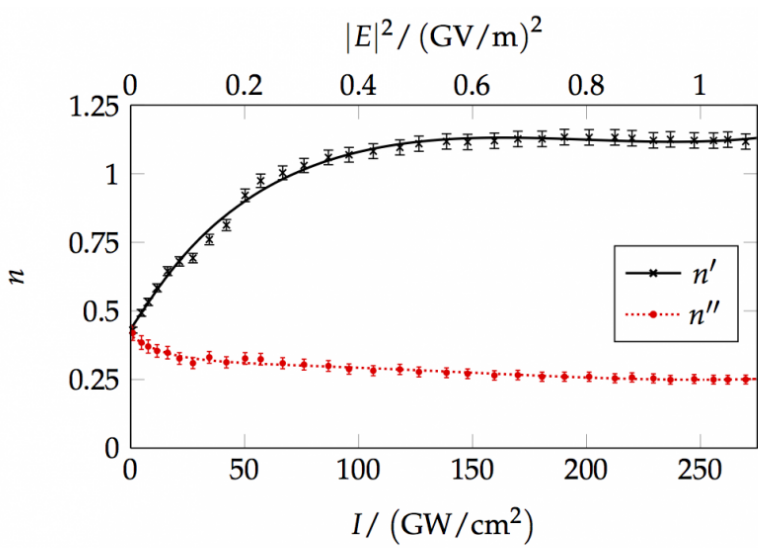
- ellipsometry



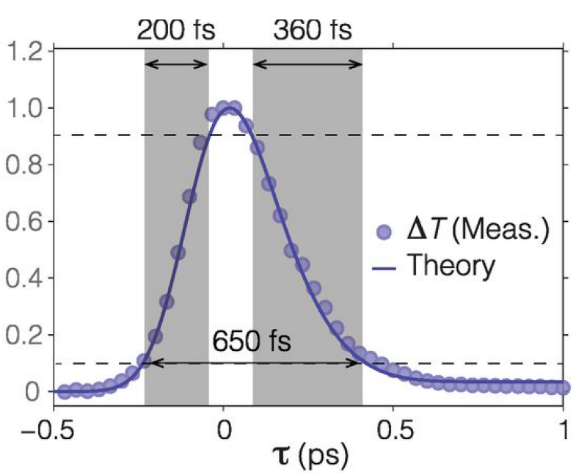
- n2 can be 3.4 x10^5 times larger than that of silica glass



- overall change in refractive index of 0.8



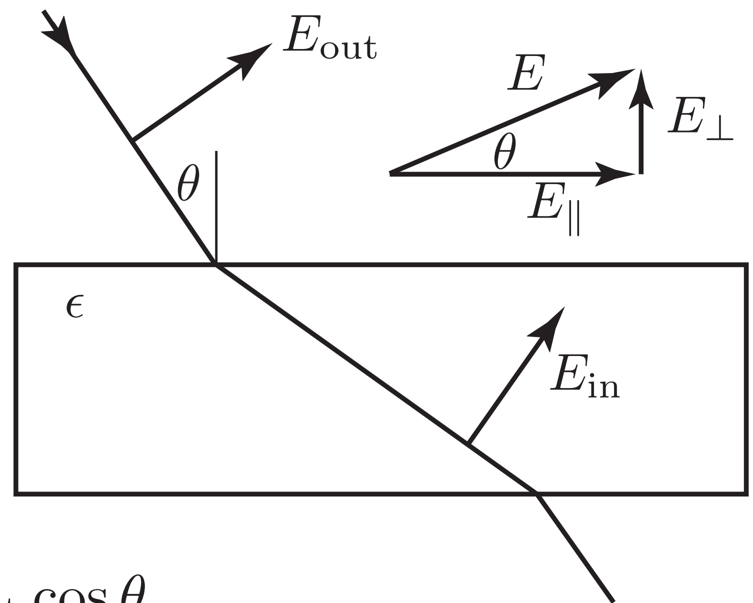
- sub picosecond reponse time



M. Z. Alam et al., Science 352, 795-797 (2016)

The NLO Response Is Larger For Oblique Incidence

We observe that the reflected wave is very weak.



Standard boundary conditions show that:

$$E_{in,\parallel} = E_{out,\parallel} = E_{out} \cos\theta$$

$$D_{in,\perp} = D_{out,\perp} \quad \Rightarrow \quad E_{in,\perp} = E_{out,\perp}/\epsilon = E_{out}\cos\theta/\epsilon$$

Thus the total field inside of the medium is given by

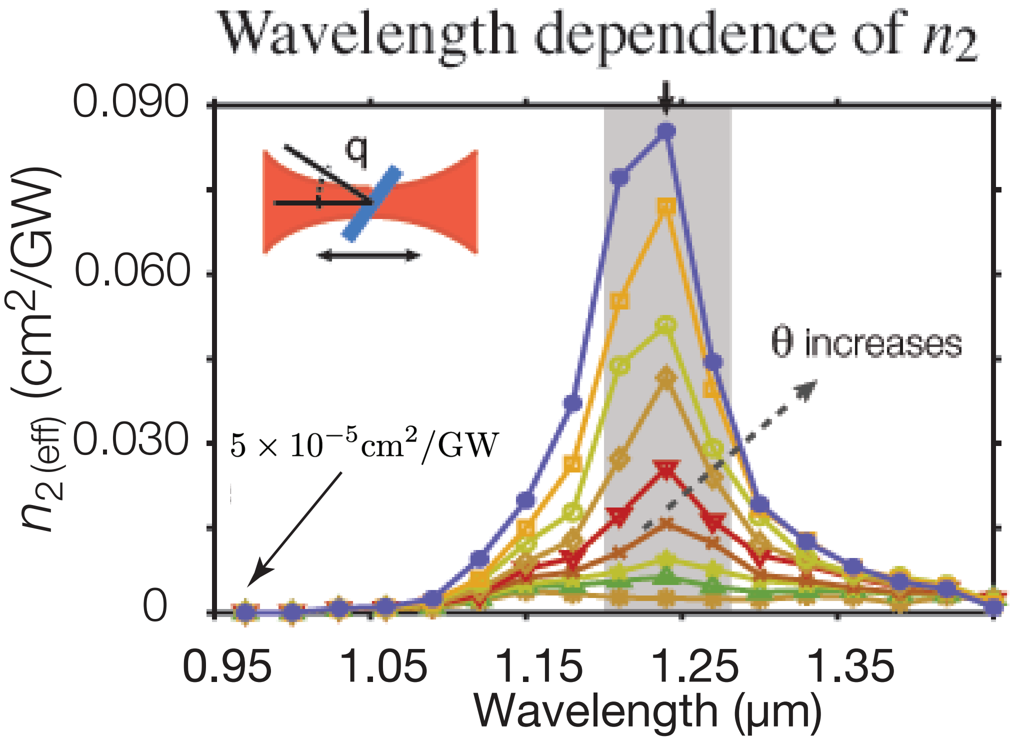
$$E_{in} = E_{out}\sqrt{\cos^2\theta + \frac{\sin^2\theta}{\epsilon}}$$

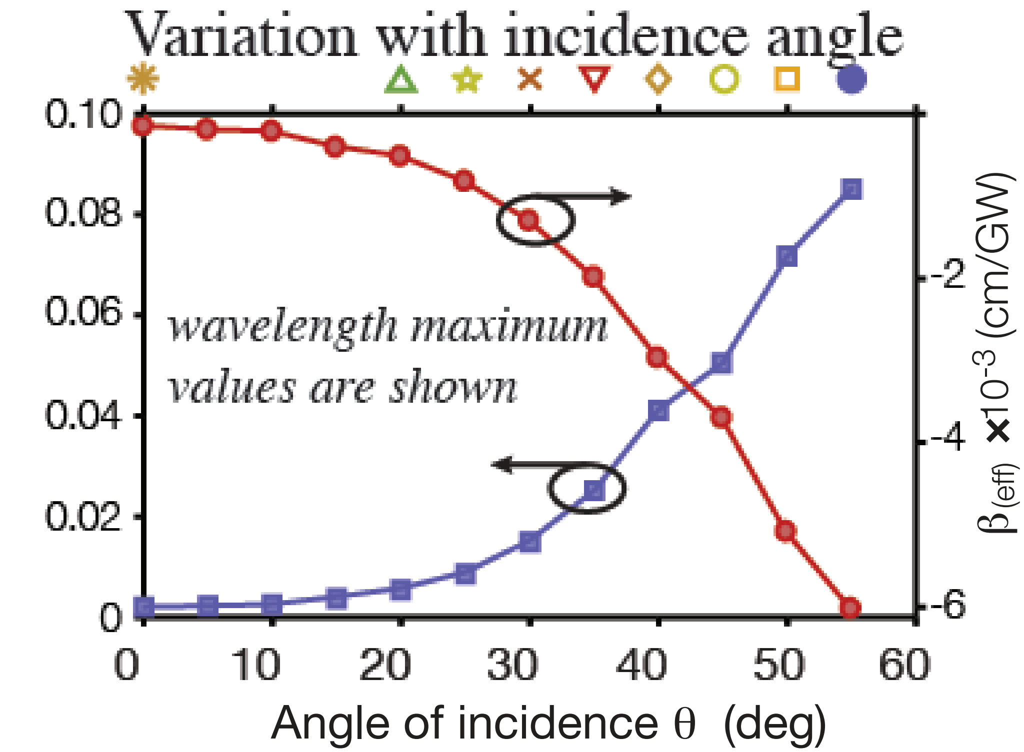
Note that, for $\epsilon < 1$, E_{in} exceeds E_{out} for $\theta \neq 0$.

Note also that, for $\epsilon < 1$, E_{in} increases as θ increases.

Huge Nonlinear Optical Response of ITO

- Z-scan measurements for various angles of incidence





peak laser intensity was 50 GW cm^{-2}

- Note that n_2 is positive (self focusing) and β is negative (saturable absorption).
- Both n_2 and nonlinear absorption increase with angle of incidence
- n_2 shows a maximum value of 0.11 cm^2/GW = 1.1 x 10^{-10} cm^2/W at 1.25 μm and 60 deg. This value is 2000 times larger than that away from ENZ region.

Why Does ENZ Lead to Large NLO Response?

1. From form of n_2 $\quad n_2 = \dfrac{3\chi^{(3)}}{4\epsilon_0 c\, n_0\, \mathrm{Re}(n_0)}$

2. From simple math: $n = n_b + \Delta n$ and $\epsilon = \epsilon_b + \Delta\epsilon$

$$n = \sqrt{\epsilon_b + \Delta\epsilon} \approx \sqrt{\epsilon_b}\left(1 + \frac{\Delta\epsilon}{2\epsilon_b}\right) = n_b + \frac{\Delta\epsilon}{2n_b}$$ and thus $$\Delta n = \frac{\Delta\epsilon}{2n_b}$$

3. Note behavior of wave equation for $\epsilon = 0$

$$\nabla \times \nabla \times \mathbf{E} + \frac{\epsilon\mu}{c^2}\frac{\partial^2}{\partial t^2}\mathbf{E} = -\mu\frac{\partial^2 \mathbf{P}^{\mathrm{NL}}}{\partial t^2}$$

4. From Maxwell's equations, it is easy to show that the nonlinear response scales as

$$\frac{\left.\dfrac{dH_x}{dz'}\right|_{\mathrm{nl}}}{|H_x|} \propto \sqrt{\frac{\mu_r}{\epsilon_r}}$$

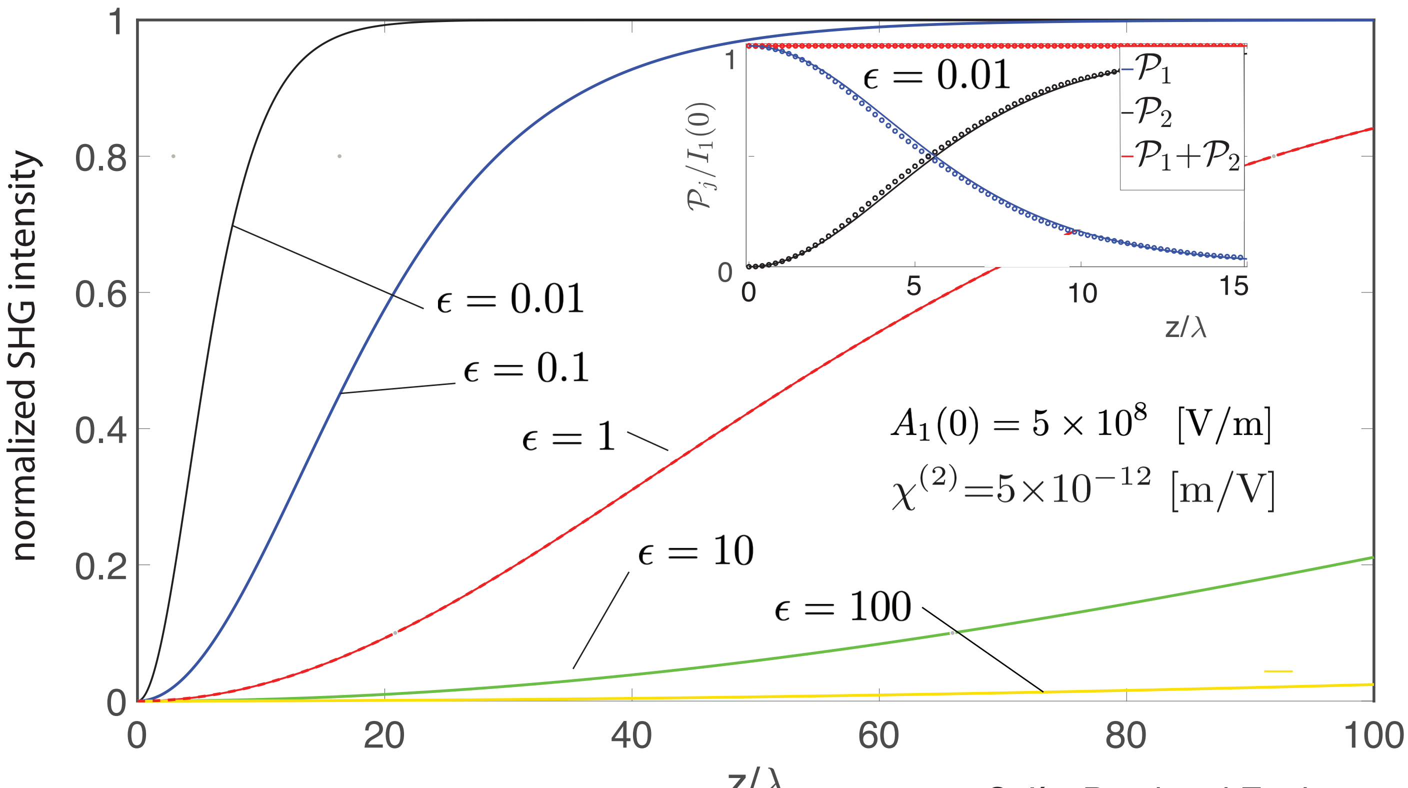
5. Detailed numerical integration confirms this behavior.

Dependence of Second-Harmonic Generation on the Linear Dielectric Permittivity

- We solve the standard equations for second-harmonic generation

$$\frac{dA_1}{dz} = i\frac{\eta_1\omega_1\chi^{(2)}}{c}A_2(z)A_1^*(z)e^{-i\Delta kz},$$
$$\frac{dA_2}{dz} = i\frac{\eta_2\omega_2\chi^{(2)}}{2c}A_1^2(z)e^{i\Delta kz},$$

- We take $\Delta k = 0$ and plot the solution for various values of the permittivity ε.
- We find that the growth rate increases dramatically as the permittivity is decreased.



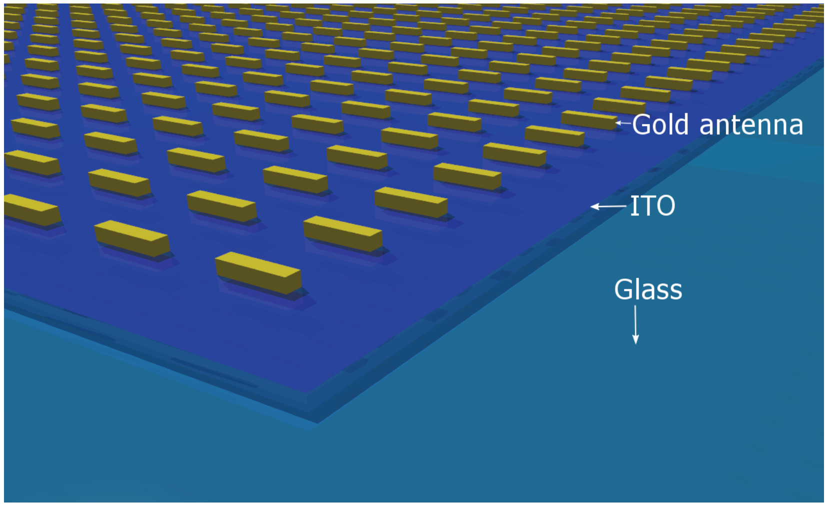
Solís, Boyd and Engheta, to be published

See also Solís, CLEO Poster JW2D.15

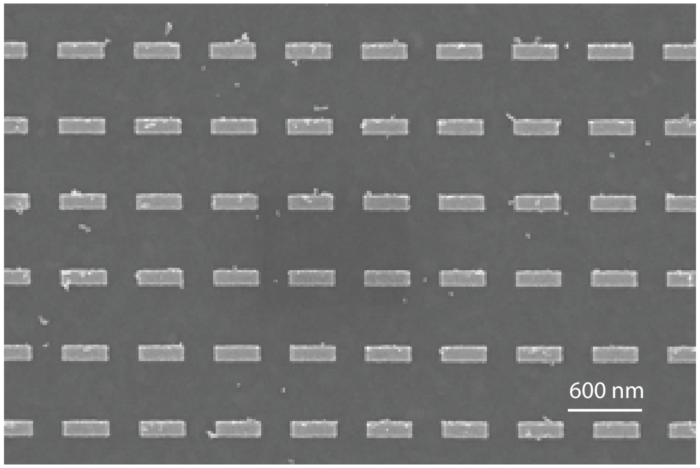
An ENZ Metasurface

- We functionalize ITO by creating a photonic metasurface
- We obtain an even larger NLO response by placing a gold antenna array on top of ITO.
  - Lightning rod effect: antennas concentrate the field within the ITO
  - Coupled resonators: ENZ resonance and nano-antennas

Concept:

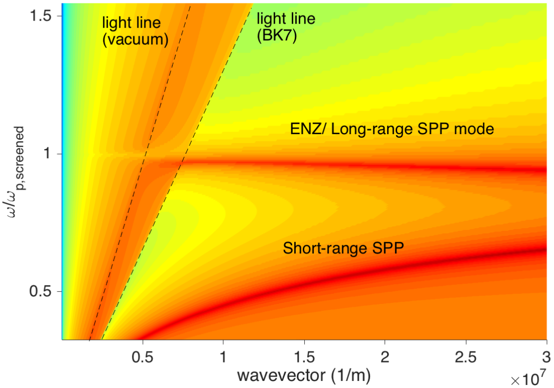


SEM:



Alam, Schulz, Upham, De Leon and Boyd, Nature Photonics 12, 79-83 (2018).

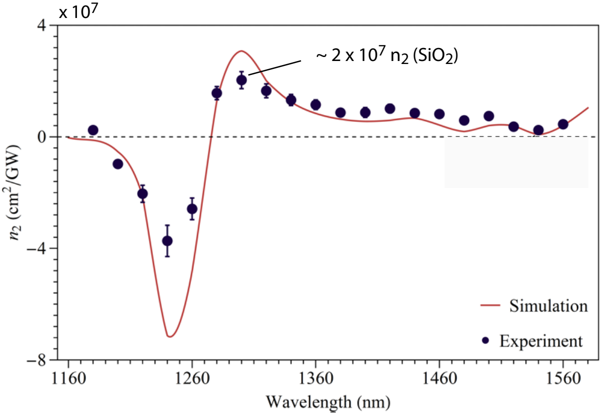
A thin ENZ medium supports a bulk plasma mode.



A thin layer of ITO supports two modes

- the bulk plasma mode, also called the ENZ or long-range SPP mode
- the short range surface plasmon polariton (SPP) mode

NLO response of the coupled antenna-ENZ system



The structure exhibits and extremely large n2 value over a broad spectral range. The on-resonance n2 value is seven orders of magnitude larger than that of silica glass.

Alam, Schulz, Upham, De Leon and Boyd,
Nature Photonics 12, 79-83 (2018).

Physics and Applications of Epsilon-Near-Zero Materials

- Physics of ENZ Materials
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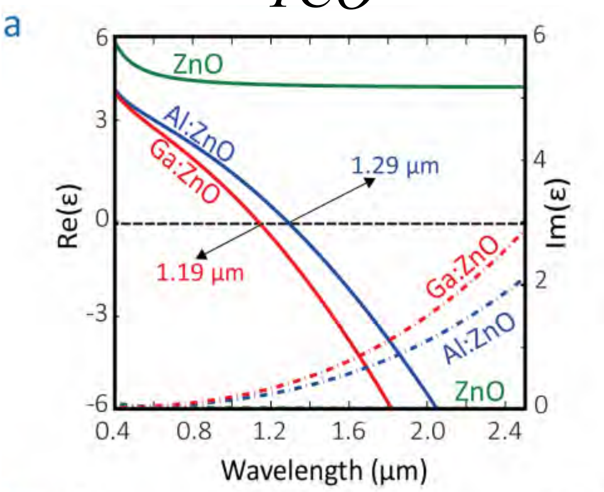
Epsilon-Near-Zero (ENZ) and Near Zero-Index (NZI) Materials

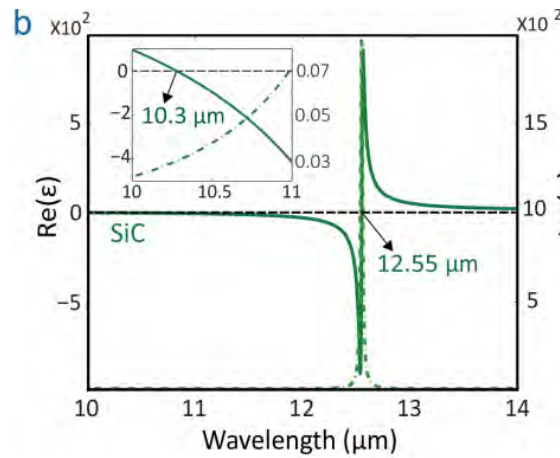
Homogeneous materials

TCO



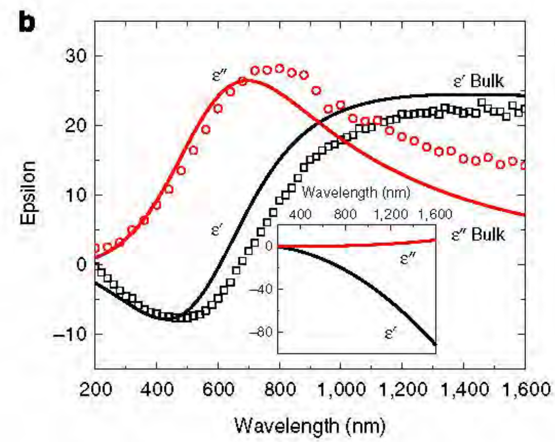
A. Boltasseva (Purdue)
Kim et al., Optica (2016)

SiC



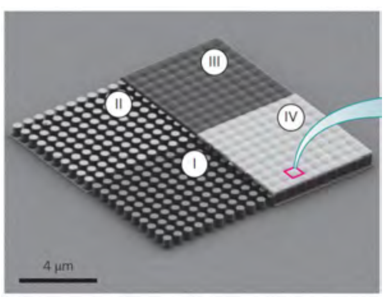
J. Caldwell (Vanderbilt)
Kim et al., Optica (2016)

$Bi_{1.5}Sb_{0.5}Te_{1.8}Se_{1.2}$

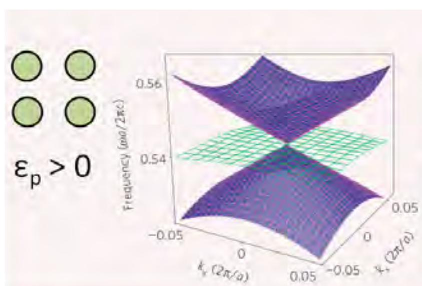


N. Zheludev (Southmapton)
Ou et al., Nat. Commun. (2014)

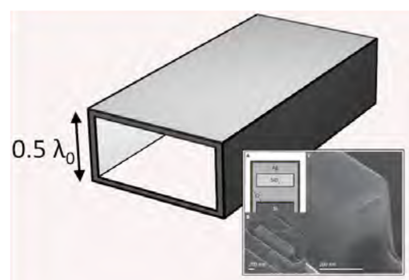
Metamaterials



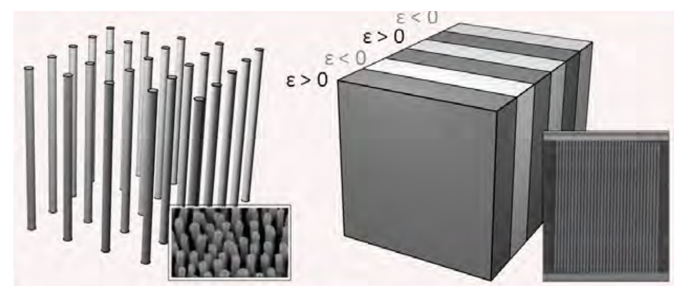
E. Mazur *Li et al., Nat. Photon.* (2015)



Chan, *Huang et al., Nat. Mater.* (2011)



SEM from: Polman's & Engheta's
Vesseur et al., PRL (2013)



$\mathrm{Re}(\varepsilon) \cong 0$

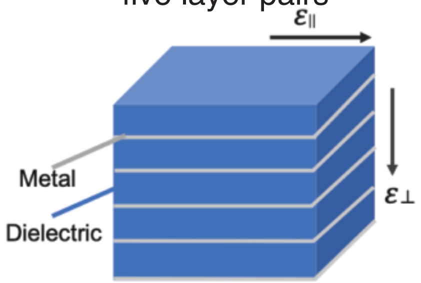
Wire SEM from: Zayat & Podolskiy
Pollard et al., PRL (2009)
StackSEM from: Polman & Engheta
Mass et al., Nat. Photon. (2013)

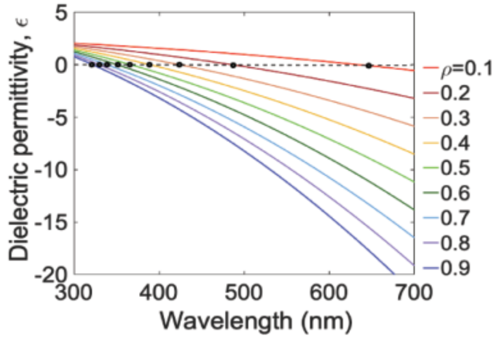
Nonlinear Optical Properties of a Layered Metamaterial in its ENZ Region

Do layered metamaterials also show enhanced NLO response at ENZ wavelength?

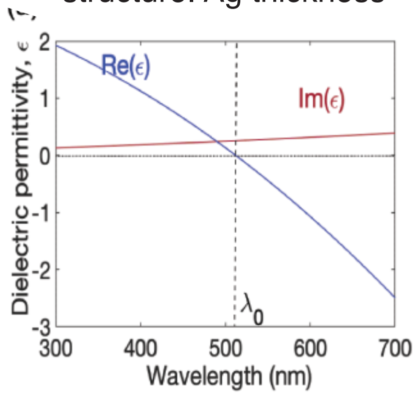
Can we use an effective-medium value of epsilon to determine the ENZ wavelength?

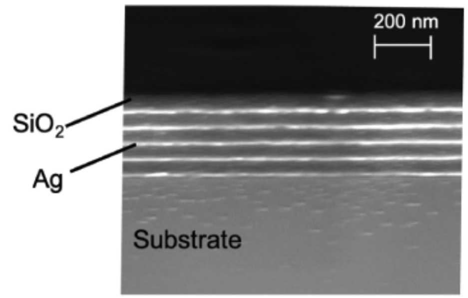
- By controlling the metallic fill fraction ρ, we can set the ENZ wavelength to be anywhere from 300 to 700 nm. We use ρ = 0.2, which corresponds to 500 nm. We deposit five layer pairs



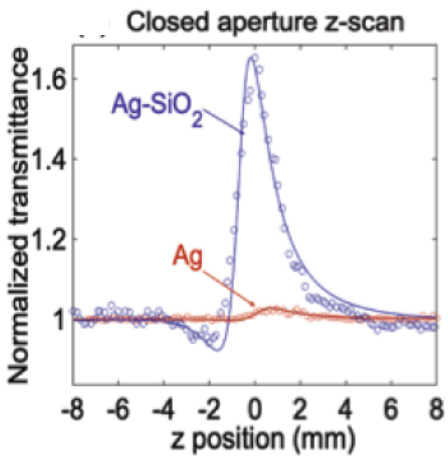


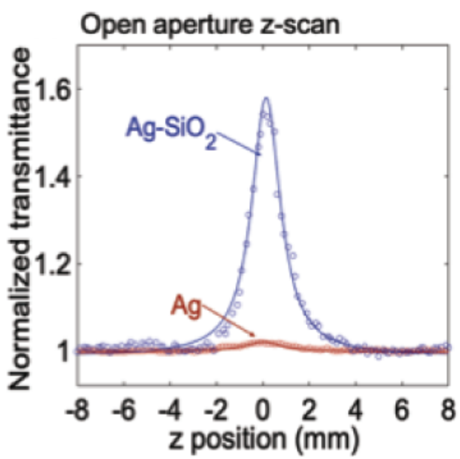
- Note that the real part of epsilon vanishes at 508 nm, close to the design wavelength. The SEM shows our structure. Ag thickness = 16 nm; silica thicness = 65 nm



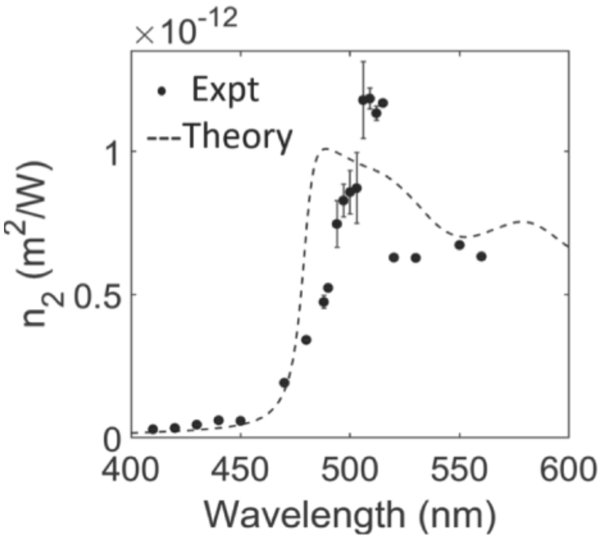


- We perform Z-scan measurements on the sample. Note the enhanced response of the composite as compared to a single layer of silver.





- Note the pronounced peak in the value of n_2 around the ENZ wavelength. We find a good but not perfect agreement with a simple effective medium theory.



Suresh, Reshef, Alam, Upham, Karimi and Boyd

Physics and Applications of Epsilon-Near-Zero Materials

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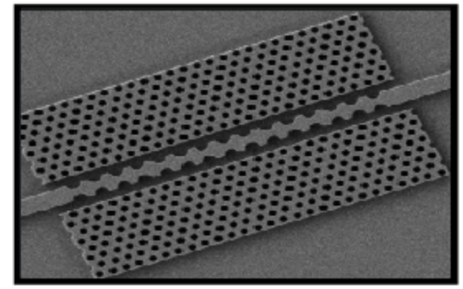
- Applications of ENZ Materials

Relaxed Phase-Matching Requirements in ENZ Media



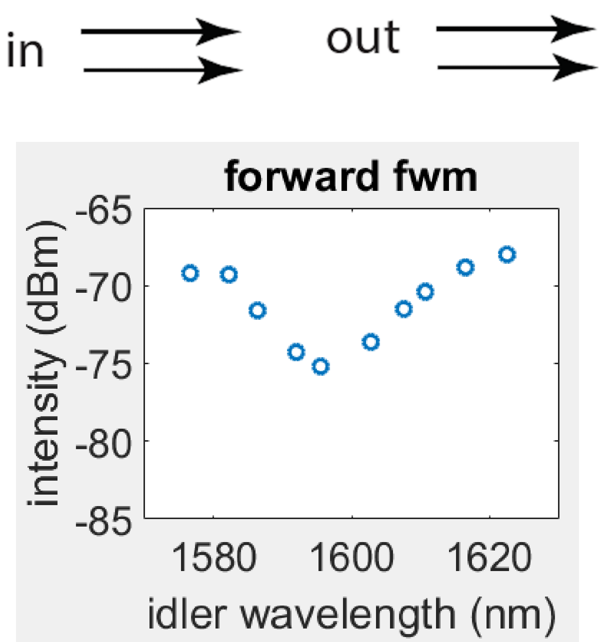
- We study four-wave mixing in a zero-index waveguide

$$2\omega_p = \omega_s + \omega_i$$

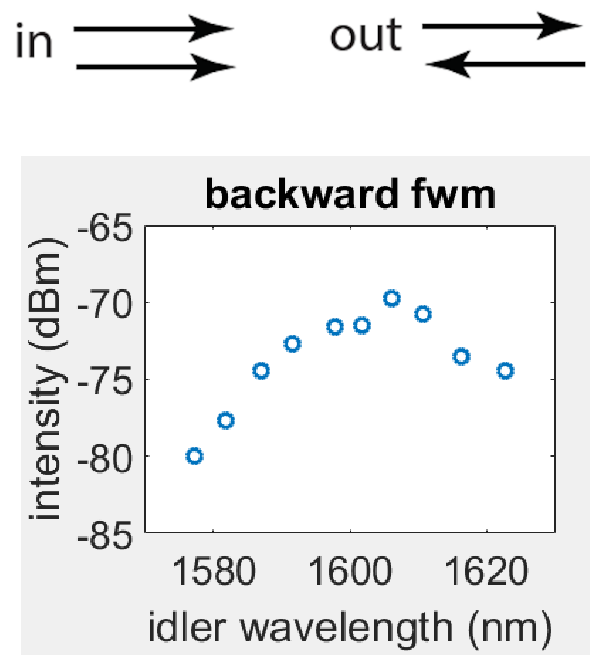
- We find that an idler field is generated in both the forward and backward directions!

- Recall that we need Δk = 0, but when n = 0, k = n ω /c vanishes for each of the interacting waves and thus so does Δk.

"forward" process



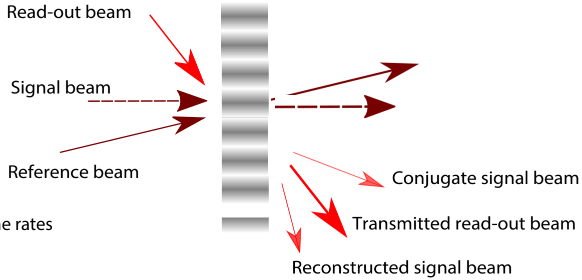
"backward" process



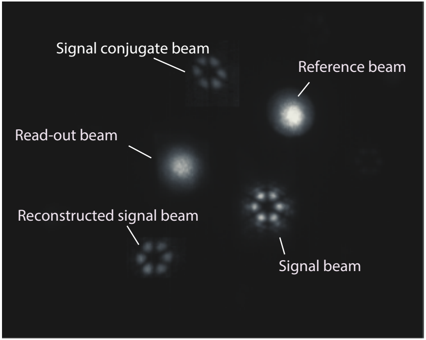
- Significance: Nonlinear optical processes that were previously believed to be too weak to be useful can be excited through use of ENZ materials.

Real-Time Holography with THz Refresh Rates

- Goal: Real-time holography with video or much faster refresh rates.
- The ultrafast response of ITO permits THz refresh rates
- Important applications involve image processing and signal processing
- Current real-time holographic materials cannot even support video frame rates



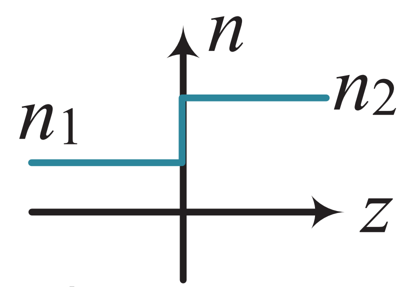
- Demonstration of image processing (edge enhancement)



Alam, Fickler, Reshef, Giese, Upham, and Boyd

Adiabatic Wavelength Conversion through Time Refraction

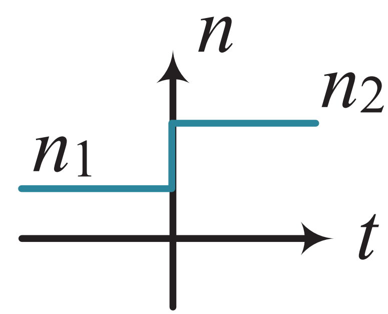
- Space refraction (e.g., Snell's law)



index vs. position

$$\frac{c}{f} = n \cdot \lambda \longrightarrow n_1\lambda_1 = n_2\lambda_2$$

- Time refraction (analog of Snell's law)



index vs. time

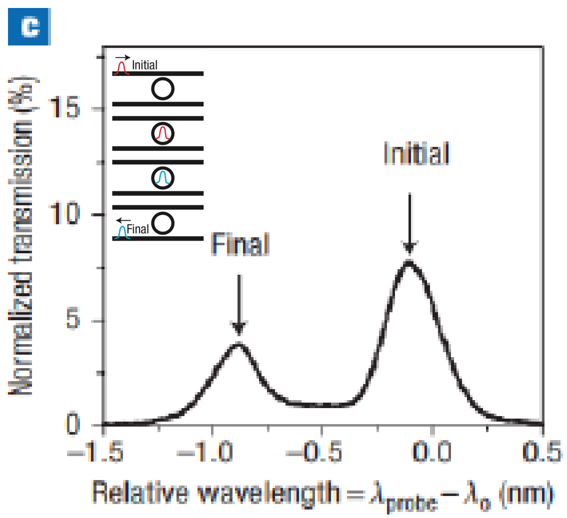
$$\frac{c}{f} = n \cdot \lambda \longrightarrow \boxed{n_1 f_1 = n_2 f_2}$$

Photon frequency (energy) is changed because of the temporal change in index, but the wavelength (inverse of momentum) is conserved in the absence of any spatial asymmetry

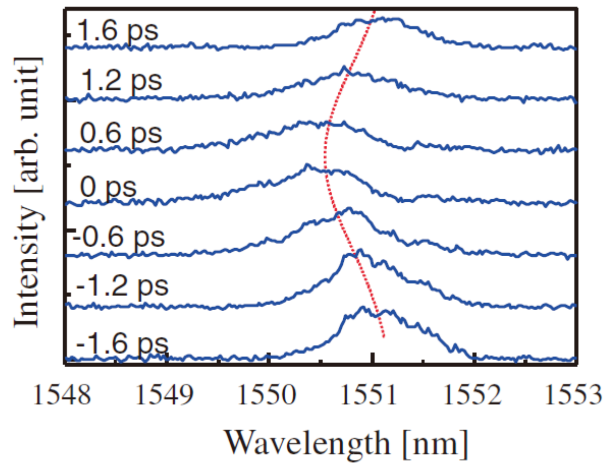
- Time refraction is an alternative way of understanding frequency broadening and shifting by self-phase modulation:

$$\delta\omega(t) = \frac{d}{dt}\phi_{\mathrm{NL}} = \frac{d}{dt}[n_2 I(t)\omega/c]$$

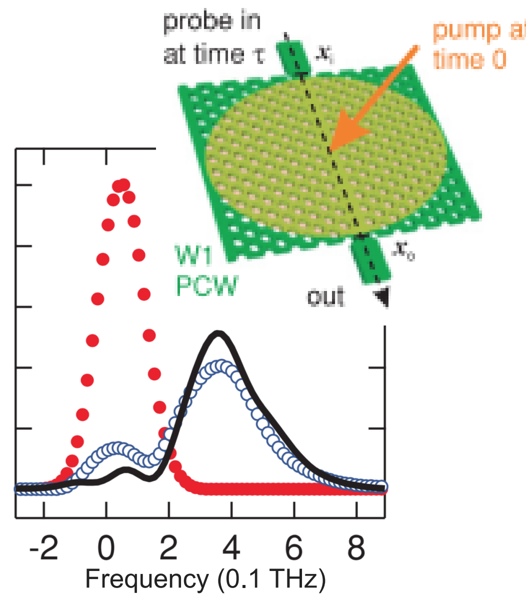
Previous work on adiabatic wavelength conversion



Nat. Photonics, 2007, 1(5): 293



Appl. Phys. Ex., 2010, 3(6): 062001.



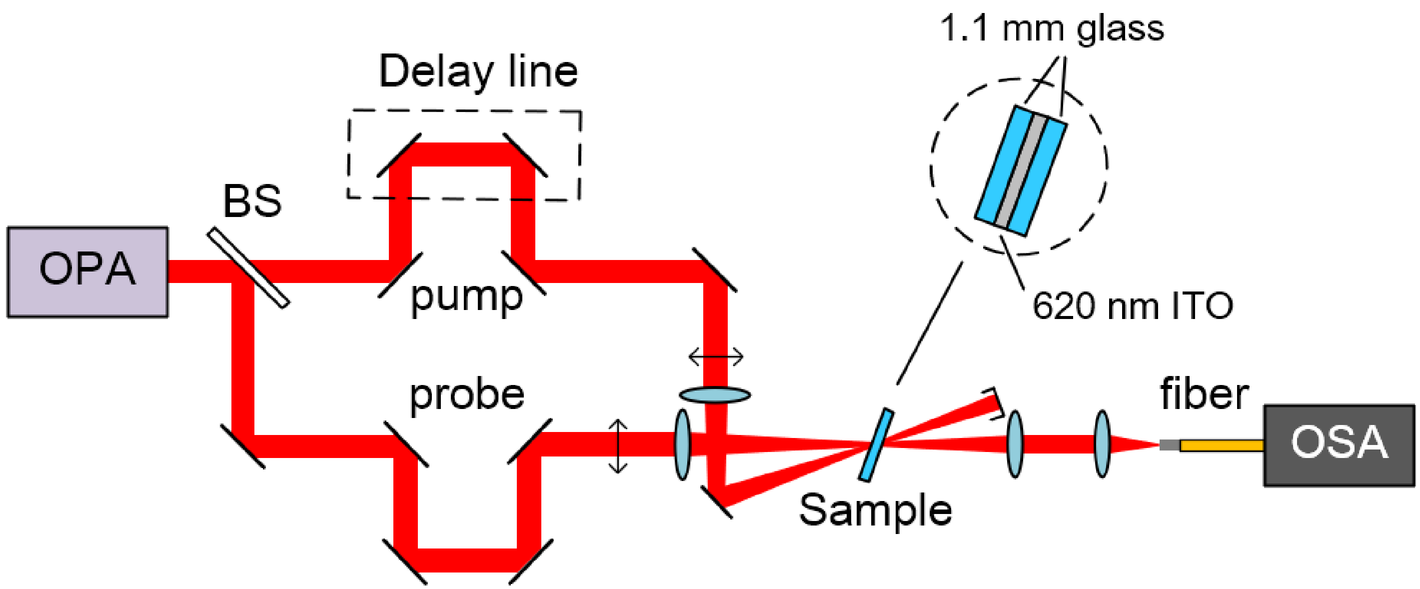
Phys. Rev. A., 2010, 81: 043837.

For small change in index $\Delta n \sim 10^{-4}$, one gets a small change in frequency (e.g., 0.06%)

But with ENZ materials we can obtain a much larger Δn

Laboratory Study of Wavelength Conversion by Time Refraction

- Pump beam creates a time-varying refractive index in ITO sample
- Frequency of probe beam is thereby modified.



OPA = optical parametric amplifier

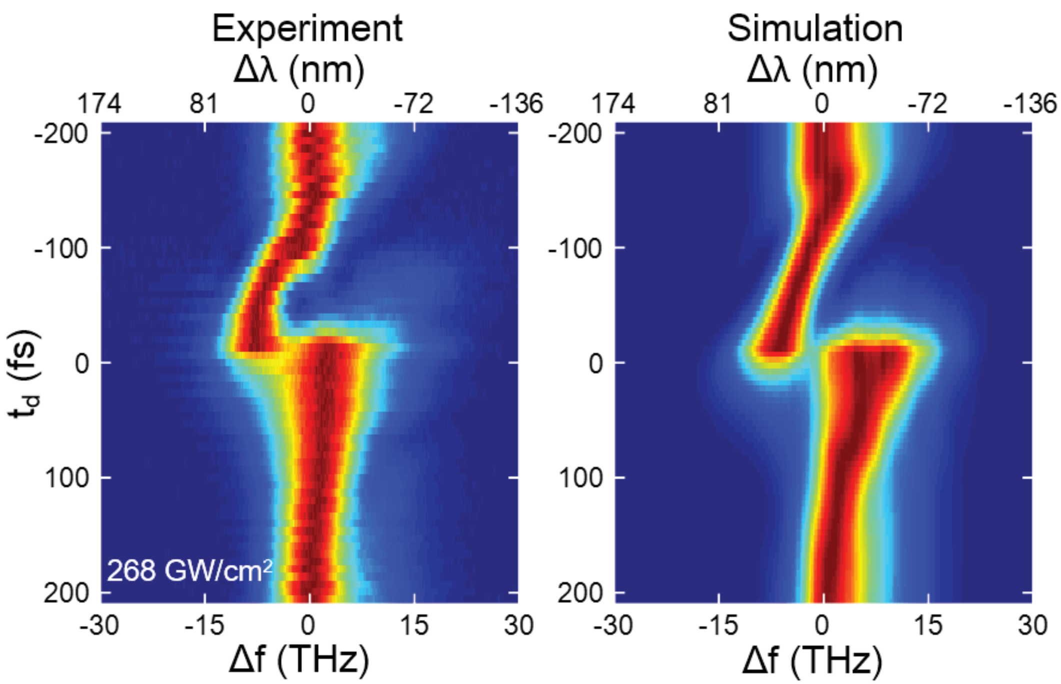
wavelength = 1240 nm

pulse duration = 120 fs

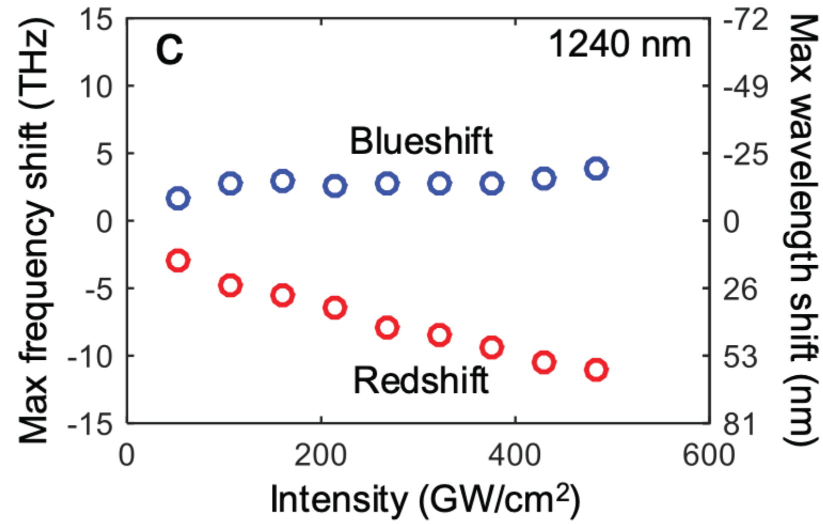
OSA = optical spectrum analyzer

Results: Adiabatic Wavelength Conversion by Time Refraction

Experimental results at 1240 nm



Probe phase and amplitude are measured by frequency-resolved optical gating (FROG)



The wavelength shift can be controlled by the pump intensity and the sign of the time delay.

- The observed effect is 100 times larger with almost 100 times smaller propagation distance than previous reports of AWC.
- Application: wavelength-division multiplexing for telecom

Zhou, Alam, Karimi, Upham, Reshef, Liu, Willner and Boyd, Nature Commun. 11:2180 (2020)

Summary: Physics and Applications of ENZ Materials

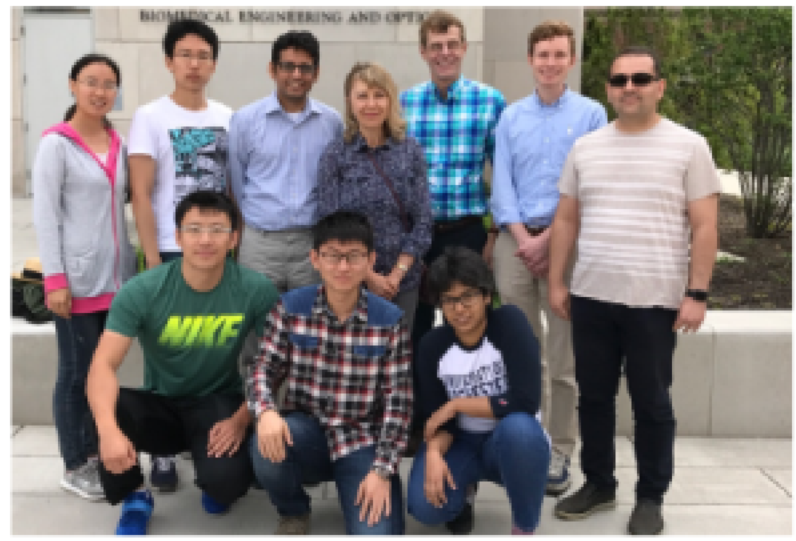
- Extremely interesting physical processes occur in ENZ materials
- ENZ materials, metamaterials, and metastructures display extremely large NLO response
- The huge, ultrafast NLO response of ENZ materials lend themselves to many important applications

The visuals of this talk are posted at boydnlo.ca/presentations

Special thanks to:

- Nader Engheta, Eric Mazur and Alan Willner for close collaboration throughout this project.
- DARPA, ARO, and NSERC for financial support.
- And especially to my students and postdocs

Rochester Group



Ottawa Group